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Theorem**

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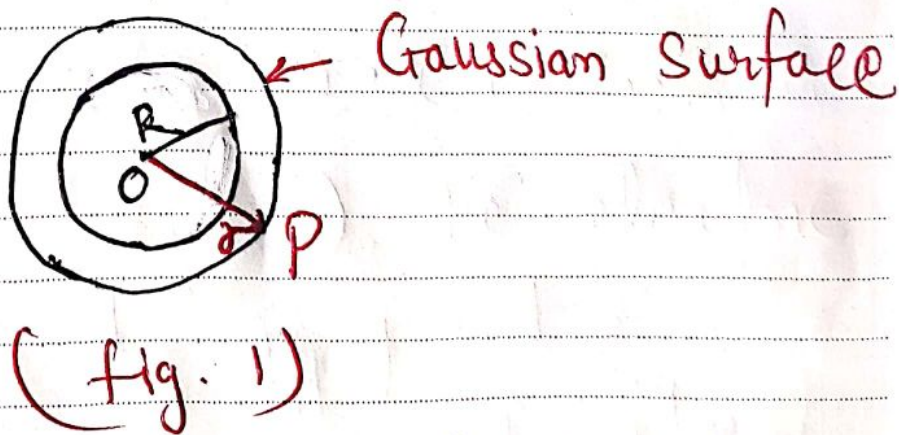
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## Applications of Gauss's law

It is possible to obtain the electric field in a simple way using the Gauss's law. This can be understood by some examples:

(1) Field outside a uniformly charged solid sphere of radius  $R$  and total charge  $q$ :

Let us draw a spherical surface at radius  $r > R$ ; this is called a "Gaussian surface".



Gauss's law says that for this surface

$$\oint_S \vec{E} \cdot d\vec{a} = \frac{q}{\epsilon_0} \quad \text{--- ①}$$



Symmetry allows us to extract  $\vec{E}$  from under the integral sign.  $\vec{E}$  points radially outward, as does  $d\vec{a}$ , so we can avoid the dot product,

$$\int_S \vec{E} \cdot d\vec{a} = \int_S |\vec{E}| da \quad \text{--- (2)}$$

and the magnitude of  $\vec{E}$  is constant over the Gaussian surface, so it comes outside the integral

$$\int_S |\vec{E}| da = |\vec{E}| \int_S da$$

( $\because$  Area of Sphere =  $4\pi r^2$ )

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$$= |\vec{E}| \cdot 4\pi r^2 \quad \text{--- (3)}$$

Thus, using eqn (1) and (3)

$$|\vec{E}| 4\pi r^2 = \frac{q}{\epsilon_0}$$

or,

$$\vec{E} = \frac{1}{4\pi\epsilon_0} \cdot \frac{q}{r^2} \cdot \hat{r}$$

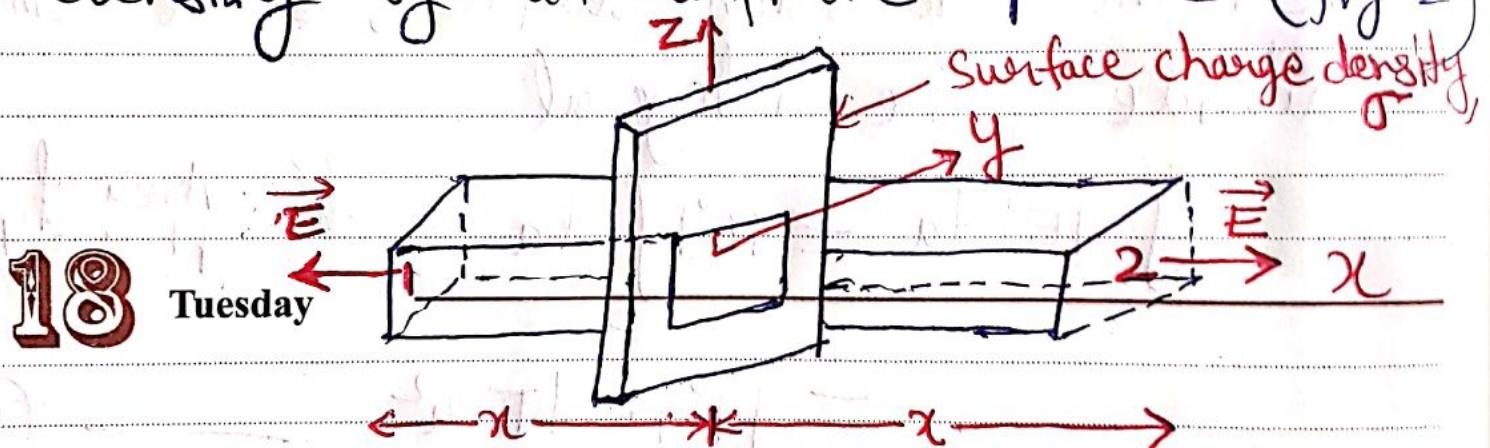
This means that the field outside the



sphere is exactly the same as it would have been if all the charge had been concentrated at the center.

## (2) Field due to a uniformly charged infinite plane sheet:

Let  $\sigma$  be the uniform surface charge density of an infinite plane. (fig. 2)



(fig. 2)  $\rightarrow$  Gaussian surface

By symmetry (from fig. 2), the electric field will not depend on  $y$  and  $z$  co-ordinates and its direction at every point must be parallel to the  $x$ -direction.

We can take the Gaussian surface to be a rectangular parallelepiped.



pipe of cross section area  $A$ .

The unit vector normal to surface 1 is in  $-x$  direction while the unit vector normal to surface 2 is in the  $+x$  direction. Therefore, flux  $\vec{E} \cdot d\vec{s}$  through both the surfaces are equal and add up. Therefore, the net flux through the Gaussian surface is  $2|\vec{E}|A$ .

or, 
$$\int \vec{E} \cdot d\vec{a} = 2A|\vec{E}| \quad \text{--- (1)}$$

The charge enclosed by the closed surface is  $\sigma A$ .

Therefore by Gauss's law -

$$\int \vec{E} \cdot d\vec{a} = \frac{\sigma \cdot A}{\epsilon_0} \quad \text{--- (2)}$$

using eq<sup>n</sup> (1) and (2),

$$2A|\vec{E}| = \frac{\sigma \cdot A}{\epsilon_0}$$

or,

$$\boxed{\vec{E} = \frac{\sigma}{2\epsilon_0} \cdot \hat{n}}$$

where,  $\hat{n}$  is a unit vector pointing away from the surface.

$\vec{E}$  is directed away from the plate if  $\sigma$  is positive and toward the plate if  $\sigma$  is negative.